

13.1 No: consider the definition of efficiency for a heat pump.

13.2  $\eta = 1 - 273/373 = 0.27$ .

13.3 Law I, Law II respectively.

13.4 Label the points A:  $(p_1, V_1, T_A)$ , B:  $(p_1, V_2, T_B)$  and C:  $(p_2, V_2, T_C)$ . The heat out on the isobar AB is  $Q_1 = C_p(T_A - T_B) = \gamma C_V(T_A - T_B)$  as you cool, while the heat in on the isochore BC is  $Q_2 = C_V(T_C - T_B)$ , and no heat is transferred on the adiabat CA. Hence using  $PV \propto T$ , the efficiency is  $\eta = W/Q_2 = 1 - Q_1/Q_2 = 1 - \gamma(T_A - T_B)/(T_C - T_B) = 1 - \gamma(p_1 V_1 - p_1 V_2)/(p_2 V_2 - p_1 V_2)$  which gives the final result.

13.5  $Q_1 = C_V(T_3 - T_2)$  and  $Q_2 = C_V(T_4 - T_1)$  and  $pV^\gamma$  is constant on an adiabat (in this case it's better to use  $TV^{\gamma-1}$  is constant).  $\eta = 1 - Q_2/Q_1$ . The result follows after some algebra.

13.6 In steady state  $\dot{Q} = Q_2$ .

The 1st law implies  $E + Q_2 = Q_1$ .

Carnot implies:  $Q_1/T_1 = Q_2/T_2$ .

Eliminate  $Q_1$  and  $Q_2$  from these equations, for example by putting the third one into the second one and yielding

$$E + Q_2 = Q_2 T_1/T_2$$

which you can use to show that

$$E = A(T_1 - T_2)(T_1/T_2 - 1) = \frac{A}{T_2}(T_1 - T_2)^2$$

This can be expanded to give a quadratic in  $T_2$ :

$$T_2^2 - (2T_1 + E/A)T_2 + T_1^2 = 0$$

which has solutions

$$T_2 = T_1 + \frac{E}{2A} \pm \sqrt{\left(\frac{E}{2A}\right)^2 + \frac{ET_1}{A}}.$$

Thus for 30% of  $E_{\max}$  you need

$$0.3E_{\max} = \frac{A}{293}10^2$$

and for 100% of  $E_{\max}$  you need

$$E_{\max} = \frac{A}{293}(\Delta T)^2$$

so that  $T_1 = 20^\circ + \Delta T = 38.3^\circ\text{C}$ .

13.7 The energy available from body 1 is  $C_p(T_1 - T_f)$ . The energy available from body 2 is  $C_p(T_2 - T_f)$ . Hence  $W = C_p(T_1 + T_2 - 2T_f)$ .

The most efficient engine is reversible and so using the Clausius theorem, the integral round a closed loop of  $dQ/T$  is zero, and hence  $\int_{T_f}^{T_1} C_p dT/T + \int_{T_1}^{T_2} C_p dT/T = 0$  and the result follows. (This result can equivalently be derived by stating  $\Delta S = 0$ , using the entropy  $S$  defined in the following chapter.)

13.8 In the steady state  $\alpha(T - T_0)$  is balanced by the heat power coming from the heat pump, call it  $Q_2 = W + Q_1$  where  $Q_1$  is the heat power extracted from the river. The efficiency

$$\eta = \frac{Q_2}{W} = \frac{T}{T - T_0},$$

so rearranging gives

$$\alpha(T - T_0) = \frac{TW}{T - T_0}$$

and hence  $TW = \alpha(T - T_0)^2$  which is a quadratic in  $T$ . Easiest perhaps to solve for  $t = T - T_0$  so that  $t^2 - tW/\alpha - T_0 W/\alpha = 0$  and the result follows (use the positive root of the quadratic or you have a cooling effect and in this country we tend to think of needing to keep our houses warmed not cooled).

# **Blundell Concepts In Thermal Physics Solutions**

**R. Shankar**



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